

SOME EXAMPLES OF ENERGETIC MATERIAL MODELLING WITH LSDYNA

Michel QUIDOT

SNPE Propulsion – Centre de Recherches du Bouchet

BP 2 – 91710 Vert -Le -Petit France

Email : m.quidot@snpe.com

ABSTRACT :

SNPE is using DYNA codes for more 15 years for characterisation and modelling of energetic materials : high explosives, solid propellants, gun propellants and pyrotechnic systems for functioning, safety, survivability and vulnerability analysis in space and defence applications.

This paper presents some examples of the use of LSDYNA in the field of solid mechanics, fluid mechanics and detonics.

The first example concerns the development of a constitutive model for a cast PBX. A general dynamic viscoelastic model developed in LSDYNA is used to analyse Split Hopkinson Pressure Bar (adapted for soft materials characterisation) experiments, reverse Taylor test instrumented by high speed framing camera and VISAR system. A postprocessing variable implemented is used for analysing dynamic failure in dynamic Brazilian tests performed with SHPB system.

Two others short examples are given : functioning of a pyrotechnic system separation in space application and LSDYNA/Euler simulation of the interaction of blast waves with explosive charges in sympathetic detonation phenomena.

MODELLING MECHANICAL BEHAVIOR OF PBX AT HIGH STRAIN RATE

Cast PBX (high explosives) are heterogeneous materials composed of a binder (a polymer such as HTPB) highly filled with energetic crystalline charges (HMX, RDX, NTO, AP, Al, ...). The general behaviour of this kind of materials is non-linear viscoelastic.

Constitutive model

We have used a generalised Maxwell viscoelastic model with n components to represent the deviatoric mechanical behaviour of cast HTPB PBX. This model was implemented in LSDYNA codes. For simplicity we use a Jaumann stress rate and an integral form of viscoelasticity which lead to the following numerical scheme :

$$s_{ij}^{n+1} = s_{ij}^n + \left(s_{ik} \omega_{kj}^{n+1/2} + s_{jk} \omega_{ki}^{n+1/2} \right) \Delta t^{n+1/2} \quad (1)$$

$$+ s_{ij}^{\nabla n+1/2} \Delta t^{n+1/2}$$

where s_{ij} is the deviatoric part of the stress tensor and ω_{ij} is the rotation tensor. The stress increment is given by :

$$\Delta s_{ij} = s_{ij}^{\nabla n+1/2} \Delta t^{n+1/2} = s_{ij}(t + \Delta t) - s_{ij}(t) \quad (2)$$

and $s_{ij}(t)$ is related to the deviatoric rate of deformation $\dot{\epsilon}'_{ij}$

$$s_{ij}(t) = \int_0^t 2G(t - \tau) \dot{\epsilon}'_{ij}(\tau) d\tau \quad (3)$$

The use of a Prony series for the shear modulus :

$$G(t) = G_0 + \sum_{i=1}^n G_i e^{-\beta_i t} \quad (4)$$

allows the calculation of the hereditary integral by recurrent formulas as often used in implicit codes. The implementation in 3D needs 5 auxiliary variables by term by using the relation $\dot{\epsilon}'_{ii} = 0$. The number of terms is not limited but only 5 are currently used. The equation of state is purely elastic. A maximum effective strain based on a Jaumann rate :

$$\epsilon_{ij}^{n+1} = \epsilon_{ij}^n + \left(\epsilon_{ik} \omega_{kj}^{n+1/2} + \epsilon_{jk} \omega_{ki}^{n+1/2} \right) \Delta t^{n+1/2} + \epsilon_{ij}^{\nabla n+1/2} \Delta t^{n+1/2} \quad (5)$$

$$\epsilon_{max} = \max_t \sqrt{2/3 \epsilon'_{ij}(t) \epsilon'_{ij}(t)} \quad (6)$$

has been added for post processing.

Experiments

Tensile static tests at different temperature and different strain rates and dynamic compression tests using Split Hopkinson Pressure Bars (SHPB) have been performed [1]. Static tests are analysed with the time-temperature superposition principle as described by Williams, Landel and Ferry [2] and values of the Young's modulus are obtained in the time range $10^{-1} - 10^8 \mu\text{s}$. A viscoelastic model is then derived in the time range of interest $0.1 - 10000 \mu\text{s}$. SHPB experiments are performed at room temperature and strain rate in the range 1000s^{-1} are achieved. Experimental stress are corrected and transformed in true (Cauchy) stress assuming incompressibility of the material.

The reverse Taylor test with Visar measurement of the rear surface is used to investigate the mechanical response of the viscoelastic material in a large continuum spectrum. Figure 1 shows the experimental arrangement.

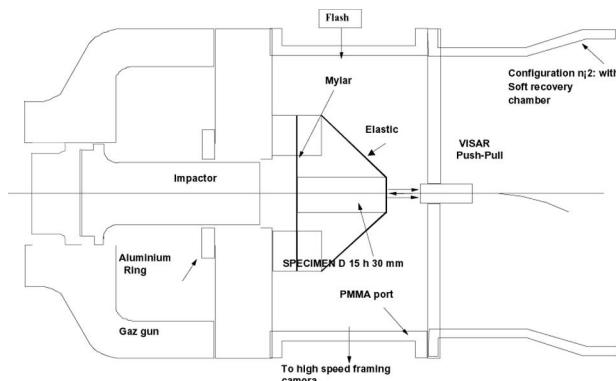


Figure 1. : Reverse Taylor Test

The tests were conducted at a low impact velocity in order to concentrate on the viscoelastic behaviour of the material before significant damage. Two tests were performed at 48 and 48.6 m/s impact velocities with well controlled impact conditions (tilt less than three mrad). In the first shot, examination of the specimen after impact shows no damage and the same dimensions as initial. The profile of the sample during impact shows a very small deformation and so the accuracy is too weak to allow a quantitative information for comparison of models. The second shot was instrumented by VISAR.

Dynamic failure properties are studied with Brazilian tests using SHPB. The deformation of the sample and propagation of failure are followed by a high speed framing camera (30000 i/s). Figure 2 shows the experimental set-up and the beginning of failure at about 1300 μs . For this test, the experimental conditions are :

- polyamide bars, diameter 40 mm, with a striker bar 1.2 m length and 9.9 m/s initial velocity
- a 36 mm diameter and 10 mm thickness sample

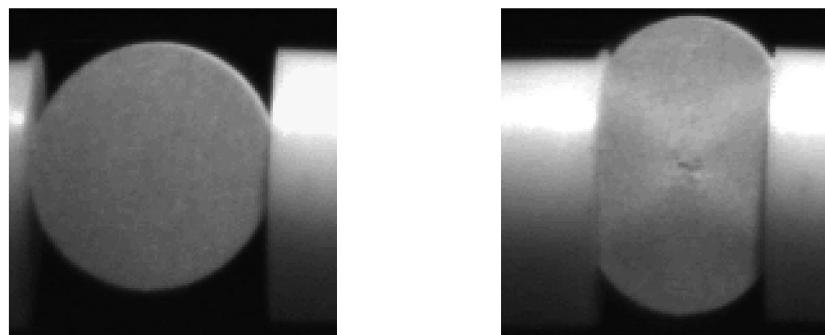


Figure 2 : Brazilian test at time 0 (left) and at first failure at the centre of the sample (right)

Numerical modelling

Figure 3 compares the VISAR experimental velocity with the numerical results with static and dynamic models. Obviously, the static model is too elastic whereas the dynamic model is quite enough dissipative but slightly too stiff at the beginning of loading. As a conclusion, the time-temperature equivalence principle does not apply for dynamic loading.

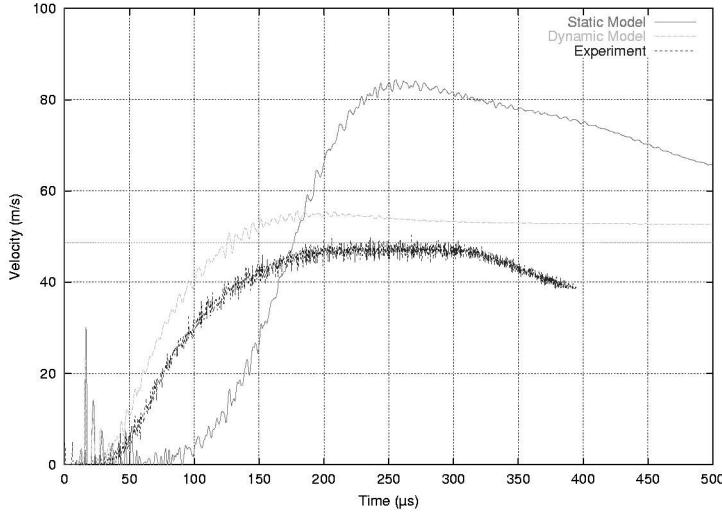
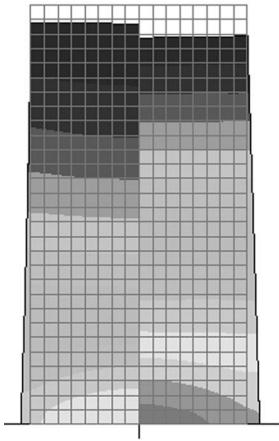


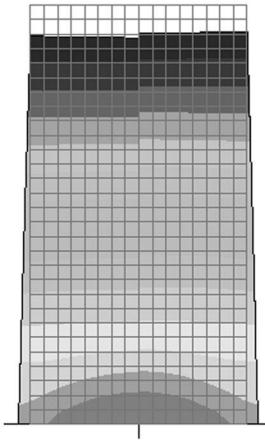
Figure 3 : Comparison of static and dynamic models on Reverse Taylor Test

In figure 4, the deformation of the sample during loading is shown, in a frame attached to the sample base, by comparison to its original size. The observed small radial deformation is consistent with the high speed framing camera result. Fringes of maximum effective strain (see equation 6) show that maximum level is achieved between 100 and 150 μ s. This maximum value of 25 % could indicate that non linear effects, not included in the viscoelastic model, exist.

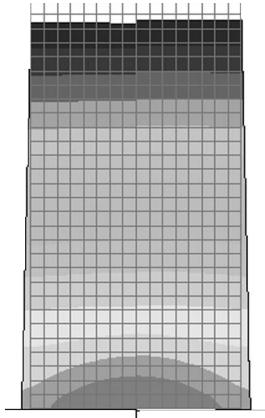
Time : 50 μ s (left) and 100 μ s (right)



Time : 150 μ s (left) and 200 μ s (right)



Time : 250 μ s (left) and 300 μ s (right)



Level



Figure 4 : Maximum effective stress in Reverse Taylor Test

For a qualitative analysis of the Brazilian test we use a simplified modelling with moving stonewalls for boundary conditions. As it is not so clear in several experiments that failure is due to shear stress or tensile stress, we implement in the viscoelastic model the postprocessing (in plane) history variable :

$$h(t)_{\max} = \underset{t}{\operatorname{Max}} \max(|\tau_{yz}(t)|, \sigma_{yy}(t)) \quad (7)$$

The numerical results (figure 5) display a stress level lower than the maximum stress (~10 MPa) measured in SHPB experiments at 750-1000 s⁻¹ but in the range of failure stress (~ 3 MPa) measured in tensile experiments performed with a high speed tensile machine at 100-250s⁻¹. The numerical simulation shows the complex loading of the sample with first localised shearing and then high tensile stress at the centre, the whole in a compression state in the z direction. This seems to be in a qualitative agreement with experiment but obviously the model is too simple to describe the complex physical phenomena in a 2D/3D stress and strain state involving large deformation.

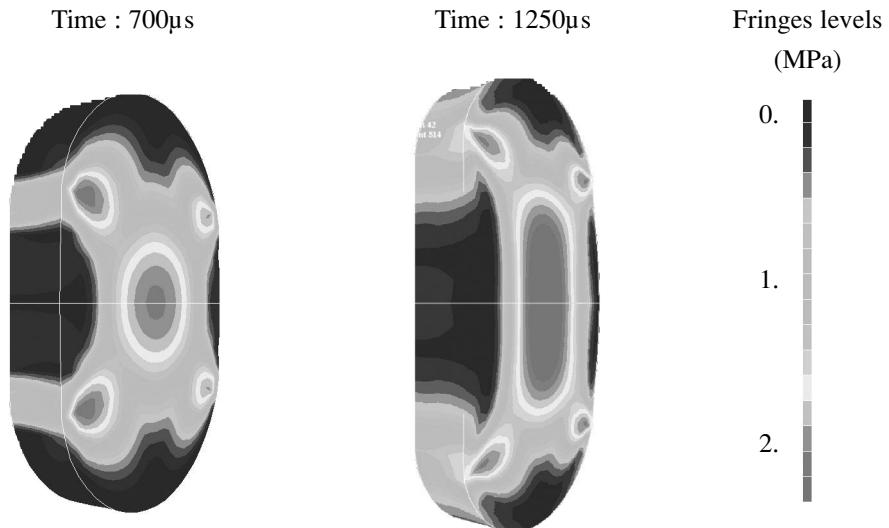


Figure 5 : Fringe of history variable h (eq. 7)

In conclusion, a viscoelastic model implemented in LSDYNA gives reasonable results for moderate loading but is insufficient for severe loading, especially for failure analysis. New developments are currently in progress as part of CMEX (Constitutive Model for Explosives) project and the first step is to separate viscoelastic effects from non linear effects.

MODELLING OF A PYROTECHNIC SYSTEM FOR SEPARATION

System studied

Pyrotechnic systems for separation are used in many space applications. A high reliability but also a low level of generated shocks in the surrounding structures are required for these systems. A numerical and experimental study has been conducted on a separation system by expanding tube for a possible optimization. This study has been supported by CNES and presented at the first European conference on Launcher Technology [5].

The functioning of this system, representative of real devices, is described figure 6. The expanding tube contains a stainless steel case, 1 mm thick, internal synthetic materials and a lead-RDX cord with 2g/m RDX charge.

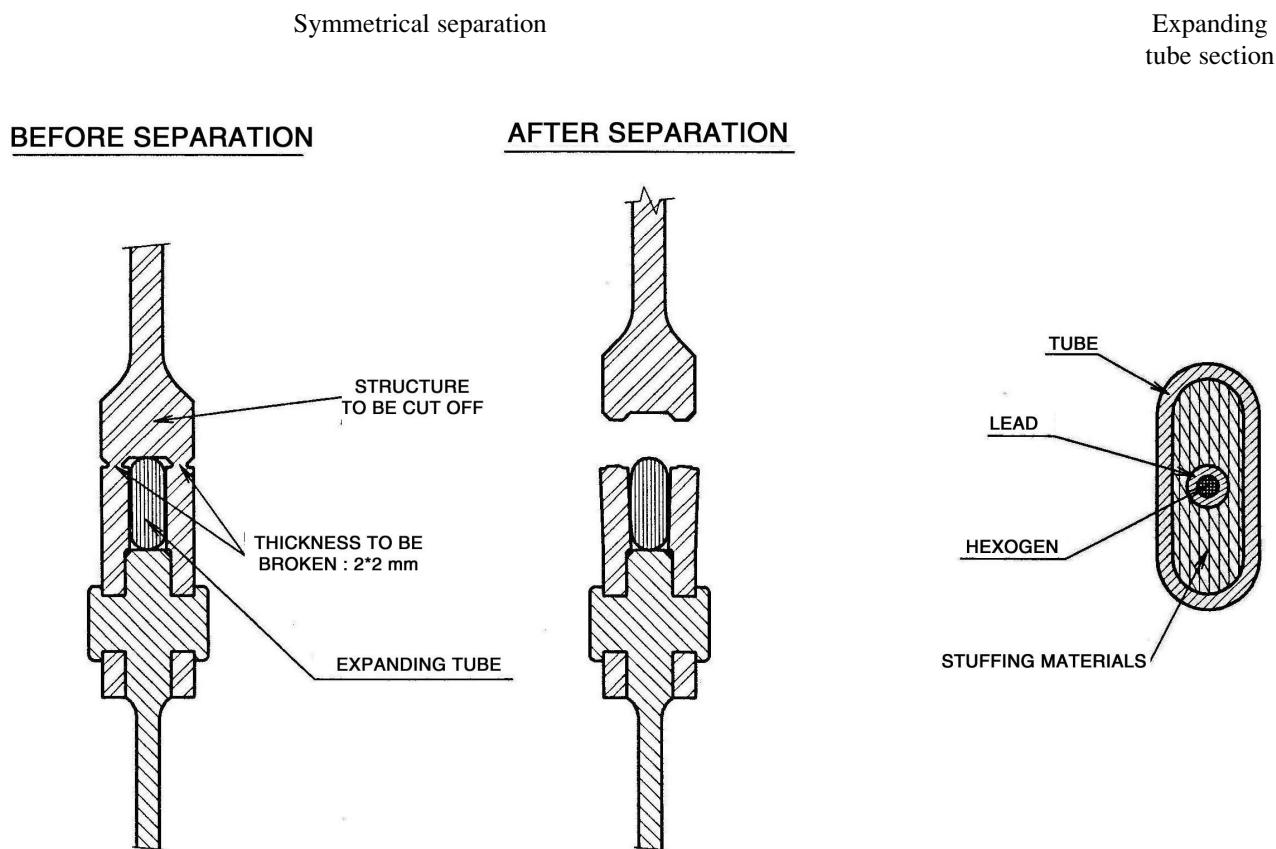
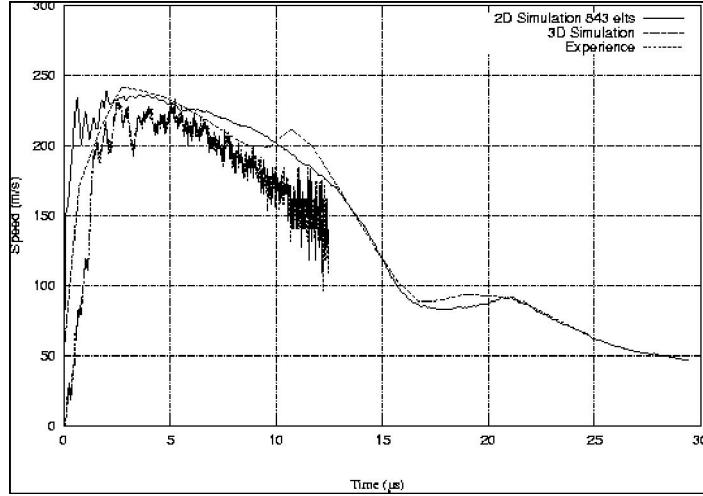


Figure 6 : Separation system

Numerical modelling

Balistic performance tests of the tube in free expansion were performed and the velocity of the minor axis (flattened form of the tube) and of the major axis (cylindrical forms of the tube) were measured by a VISAR and a streak camera [3]

Numerical simulations in 2D of the free expansion of the tube show that the two principal components are the steel case and the explosive. The elastoplastic model for the steel case is based on (static) tensile test. A first estimate of the JWL EOS of RDX at density 1.55 g/cm³ is determined by a thermochemical code [4]. Comparison with literature data [6],[7] shows that the detonation of the RDX charge is non ideal due the dimension of the explosive which is about the dimension of the reactive zone of RDX. The JWL EOS was so recalibrated [3] by a simplified energetic method based on the experimental measurements. A 3D calculation [5] demonstrates that there is no 3D effects. Figure 7 shows the good agreement between 2D and 3D calculations and experimental result. Figure 8 illustrates the functioning of the expanding tube.



**Figure 7 : Velocity at the minor axis of the tube in free expansion
Comparison of 2D and 3D simulations and experiment**

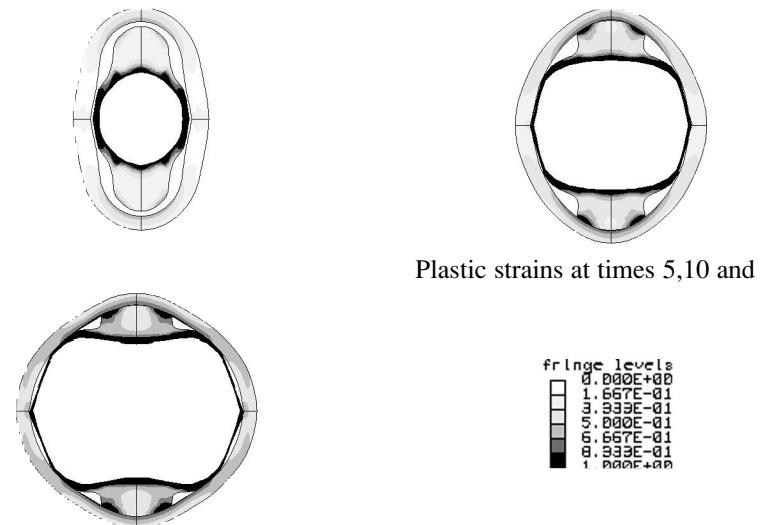


Figure 8 : Expanding tube distortions

BLAST WAVES AND SYMPATHETIC DETONATION MODELLING

LSDYNA/Euler capabilities (version 946) have been tested on sympathetic detonation phenomena : a donor charge is detonated and the reaction of an acceptor charge by the blast wave generated is observed as a function of the distance between donor and acceptor charges.

The simulation in 2D plane uses Euler, Lagrange and SALE formulations :

- the acceptor charge is Lagrangian
- the donor charge (and part of surrounding air) is first Lagrangian ($t=0. - 5.17\mu s$), then ALE ($t=5.17\mu s - 6.17\mu s$) and finally Eulerian
- the air is Eulerian (Van Leer scheme)

The explosive is COMP B and is modeled with a Lee-Tarver reaction rate and JWL EOS for both reacted and unreacted explosive determined by M. J. Murphy and all [8].

The reaction rate :

$$\frac{dF}{dt} = I(1-F)^{\frac{2}{9}}(\rho/\rho_0 - 1)^4 + G(1-F)^{\frac{2}{9}}F^{\frac{2}{3}}P^z$$

is function of compression rate (first term) and pressure (second term). The first term that represents the ignition is cut off when the fraction reacted F reaches F_{igmax} .

The JWL EOS relates pressure to specific volume and temperature :

$$P = A \exp(-R_1 V) + B \exp(-R_2 V) + \omega c_v T / V$$

The coefficients are summarized table 1.

Unreacted Explosive	
ρ_0 (g/cm3)	1.71
T_0 (K)	298
C_{vs} (Mbar*cm3 /g /K)	$2.487 \cdot 10^{-5}$
A_s (Mbar)	778.1
B_s (Mbar)	-0.05031
R_{1s}	11.3
R_{2s}	1.3
ω_s	0.8938
Detonation Products	
D_{cj} (cm/ μ s)	0.7980
P_{cj} (Mbar)	0.295
A (Mbar)	5.242
B (Mbar)	0.07678
R_1	4.2
R_2	1.1
ω	0.34
E_0 (Mbar*cm ³ /cm ³)	0.085
C_v (Mbar*cm3 /g /K)	$0.5834 \cdot 10^{-5}$
Lee-Tarver reaction rate	
I (μ s ⁻¹)	44.
$F_{igm\max}$	0.3
G (Mbar ^{-z} μ s ⁻¹)	514
Z	2.

Table 1 : Coefficients for the Lee-Tarver COMP B model

Figure 9 shows a non reaction of the acceptor for a one caliber distance between charges and figure 10 shows the detonation of the acceptor for a half caliber distance between charges. These results are in agreement with experiments performed at SNPE [9] and are encouraging for using new capabilities of LSDYNA/Euler

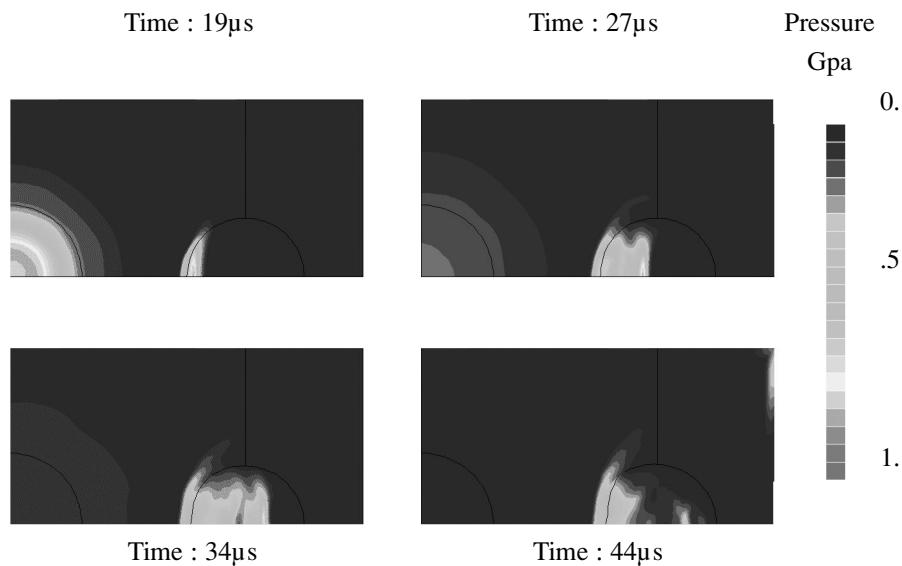


Figure 9 : Sympathetic Reaction – Donor/Acceptor distance : 1 caliber

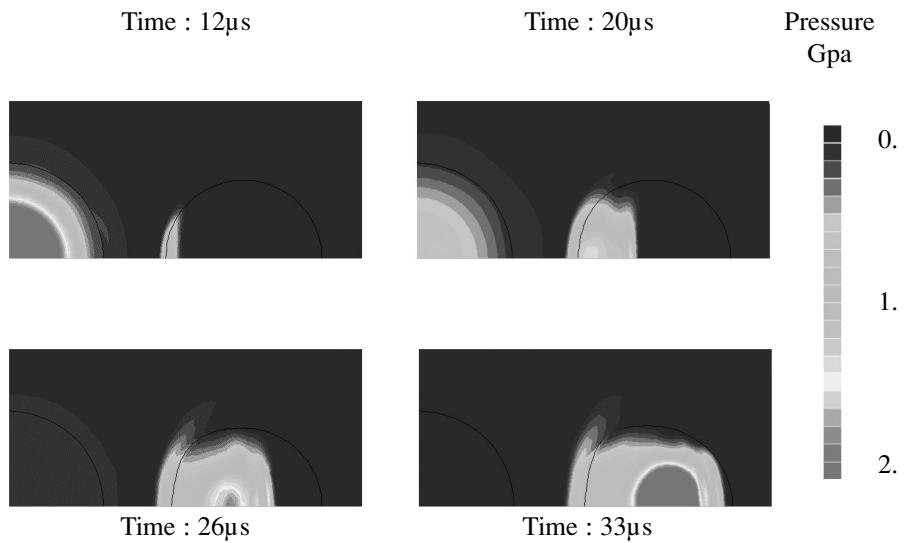


Figure 10 : Sympathetic Reaction – Donor/Acceptor distance : 1/2 caliber

CONCLUSIONS

The examples presented show the powerful capabilities of LSDYNA for characterization and modelling of energetic materials in space and defense applications. But for quantitative analysis, the accuracy of the numerical results are proportional to the materials input data. The complexity of the physics of energetic materials needs detailed models and appropriated experiments to validate the models developed.

References

- [1] M. Quidot, P. Racimor, P. Chabin
Constitutive models for cast PBX at high strain rate
Shock Compression of Condensed Matter. APS Conference 1999
- [2] Williams M.L., Landel R.F., and Ferry J.D.,
J. Amer. Chem. Soc. **77**,3701 (1955)
- [3] M. Quidot,, S. Lécume, J.P. Delon
Modélisation du fonctionnement d'un système de découpe
7^e Congrès International de Pyrotechnie 1999
- [4] F. L. Fried, W. M. Howard, P. C. Souers
Cheetah 2.0, User's Manual
- [5] M. Quidot,, S. Lécume, D. Desailly, J.P. Delon
Computational simulation of an expanding tube separation device
First European Conference on Launcher Technology 1999
- [6] B. M. Dobratz, P. C. Crawford
« LLNL Explosives handbook », 1985
- [7] T. R. Gibbs, A. Popolato
« LASL Explosive properties data »
- [8] M. J. Murphy, E. L. Lee, A. M. Weston, A. E. Williams
Modeling shock initiation in composition B
Tenth International Symposium on Detonation 1993
- [9] R. Heyzen, P. Babault, L. Allain
Rapport SNPE N°08/93/CRB/S/TS 1993