

# Determining the Material Constants for Mullins Effect in Rubber Part One: Uniaxial

William W. Feng  
John O. Hallquist

*Livermore Software Technology Corporation  
7374 Las Positas Road  
Livermore, CA 94551*

## Introduction

In this paper, the strain-energy density with Mullins damage function on unloading and subsequent reloading is considered. We introduce a damage function that has four material constants: two for unloading and two for subsequent reloading. The effect of these constants on unloading and subsequent reloading is studied for uniaxial extension. We determine these four material constants from a set of numerically generated uniaxial extension test data. The mathematical formulation has been implemented in LS-DYNA<sup>®</sup> for user application and evaluation.

This paper will be extended to two-dimensional problems and a set of biaxial test data will be obtained and analyzed. The second part of this paper will be presented in another LS-DYNA conference.

## Formulation

The strain-energy density function with Mullins damage function of a rubber is  $\tilde{W}(\lambda_i)$ , and

$$\tilde{W}(\lambda_i) = \eta W(\lambda_i) \quad (1)$$

where  $W$  is the strain-energy density function based on the initial loading, and  $\eta = \eta(W)$  is a damage function for the Mullins effect.

$$\frac{\partial \tilde{W}}{\partial \lambda_i} = \eta \frac{\partial W}{\partial \lambda_i} + W \frac{\partial \eta}{\partial \lambda_i} = \left( \eta + W \frac{\partial \eta}{\partial W} \right) \frac{\partial W}{\partial \lambda_i} \quad (2)$$

The Cauchy stresses (force per unit deformed area) are

$$t_1 = \frac{1}{\lambda_2 \lambda_3} \frac{\partial \tilde{W}}{\partial \lambda_1} \quad (3)$$

There are two similar equations for  $t_2$  and  $t_3$ .

The following damage function, a Cauchy first-order ordinary-differential equation, is chosen for this report.

$$\text{For initial loading} \quad W \frac{\partial \eta}{\partial W} + \eta = 1 \quad (4a)$$

$$\text{For unloading} \quad W \frac{\partial \eta}{\partial W} + \eta = 1 - \frac{1}{r_1} \tanh \left[ \frac{1}{m_1} \left( 1 - \frac{W}{W_m} \right) \right] \quad (4b)$$

$$\text{For subsequent reloading} \quad W \frac{\partial \eta}{\partial W} + \eta = 1 - \frac{1}{r_2} \tanh \left[ \frac{1}{m_2} \left( 1 - \frac{W}{W_m} \right) \right] \quad (4c)$$

$W_m(\lambda_i)$  is the maximum strain-energy density function before unloading.  $r_1$ ,  $r_2$ ,  $m_1$  and  $m_2$  are the material constants for the Mullins damage function. With this damage function the unloading and subsequent reloading follow different paths, as shown in Figure 1. For a loading with a value of the strain-energy density function greater than  $W_m(\lambda_i)$ , the process repeats.

For Mooney-Rivlin materials the strain-energy density equation is:

$$W = C_1(I_1 - 3) + C_2(I_2 - 3) = C_1[(I_1 - 3) + \alpha(I_2 - 3)] \quad (5)$$

where  $C_1$  and  $C_2$  are material constants and  $\alpha = C_2/C_1$ . The strain invariants  $I_1$  and  $I_2$  are written in terms of the principal stretch ratios  $\lambda_1$ ,  $\lambda_2$  and  $\lambda_3$

$$\begin{aligned} I_1 &= \lambda_1^2 + \lambda_2^2 + \lambda_3^2 \\ I_2 &= \lambda_1^2 \lambda_2^2 + \lambda_2^2 \lambda_3^2 + \lambda_3^2 \lambda_1^2 \end{aligned} \quad (6)$$

An incompressibility condition is assumed in the Mooney-Rivlin material constitutive equation, so that

$$\lambda_1 \lambda_2 \lambda_3 = 1 \quad (7)$$

For uniaxial tension or compression in 1-direction,  $\lambda_2 = \lambda_3$  and  $\lambda_2^2 = \lambda_1^{-1}$

Hence,

$$\begin{aligned} I_1 &= \lambda_1^2 + 2/\lambda_1 \\ I_2 &= 2\lambda_1 + 1/\lambda_1^2 \end{aligned} \quad (8)$$

Cauchy stress  $t_1$  is related to the uniaxial stretch ratio  $\lambda_1$ .

$$\text{For initial loading} \quad t_1 = 2C_1 \left( \lambda_1^2 - \frac{1}{\lambda_1} \right) \left( 1 + \frac{\alpha}{\lambda_1} \right) \quad (9a)$$

$$\text{For unloading} \quad t_1 = 2C_1 \left\{ 1 - \frac{1}{r_1} \tanh \left[ \frac{1}{m_1} \left( 1 - \frac{W}{W_m} \right) \right] \right\} \left( \lambda_1^2 - \frac{1}{\lambda_1} \right) \left( 1 + \frac{\alpha}{\lambda_1} \right) \quad (9b)$$

$$\text{For subsequent reloading} \quad t_1 = 2C_1 \left\{ 1 - \frac{1}{r_2} \tanh \left[ \frac{1}{m_2} \left( 1 - \frac{W}{W_m} \right) \right] \right\} \left( \lambda_1^2 - \frac{1}{\lambda_1} \right) \left( 1 + \frac{\alpha}{\lambda_1} \right) \quad (9c)$$

The result for uniaxial extension obtained from EXCEL calculations is shown in Figure 1. The material constants are:

$$C_1 = 50, \alpha = 0.1, r_1 = 0.8, m_1 = 1.0, r_2 = 0.5 \text{ and } m_2 = 5.$$

The material is first stretched to point (1) followed by the initial loading path. After point (1) the material is unloaded to the unstretched state and then reloaded to point (1) again. The stress-stretch ratio curves are shown in the figure and the Mullins effect is clearly seen. The material is then further stretched past point (1) to point (2). No Mullins effect occurs after point (1). The material is then unloaded again to the origin and reloaded again to point (2). The Mullins effect is seen again.

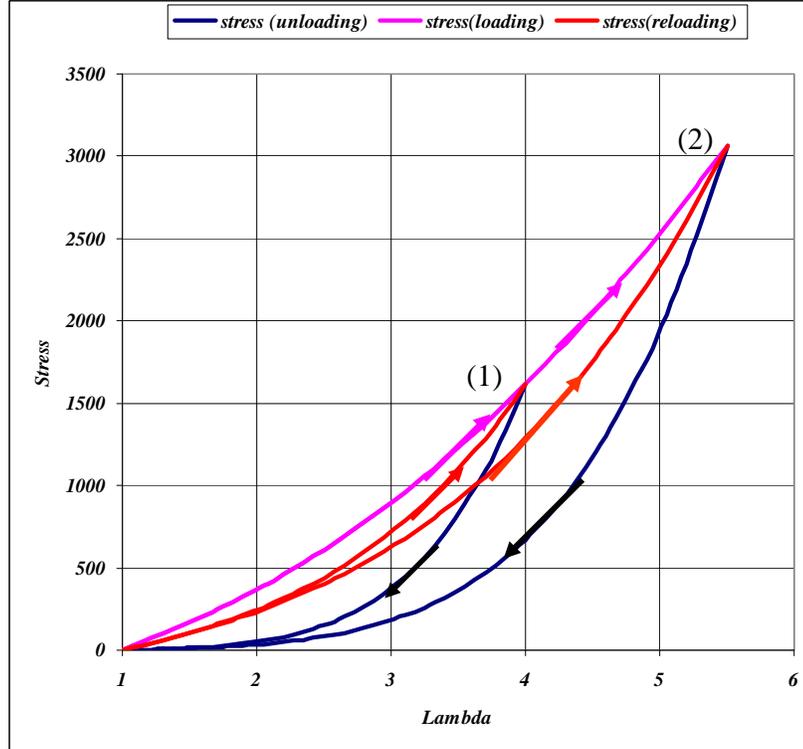


Figure 1. The Mullins effect for uniaxial loading, unloading and subsequent reloading, obtained from EXCEL.

**The effects of constants  $r_1$  and  $m_1$**

The effect of constants  $r_1$  and  $m_1$  on unloading is studied for uniaxial extension. The effect of  $r_1$  is shown in Figure 2. The effect of  $m_1$  is shown in Figure 3. They cover a wide range of unloading values. When either  $r_1$  or  $m_1$  approaches a large value, then

$$\frac{1}{r_1} \tanh\left[\frac{1}{m_1}\left(1 - \frac{W}{W_m}\right)\right] \rightarrow 0, \text{ and } W \frac{\partial \eta}{\partial W} + \eta \rightarrow 1. \tag{10}$$

The loading and unloading curves coincide and the Mullins effect vanishes.

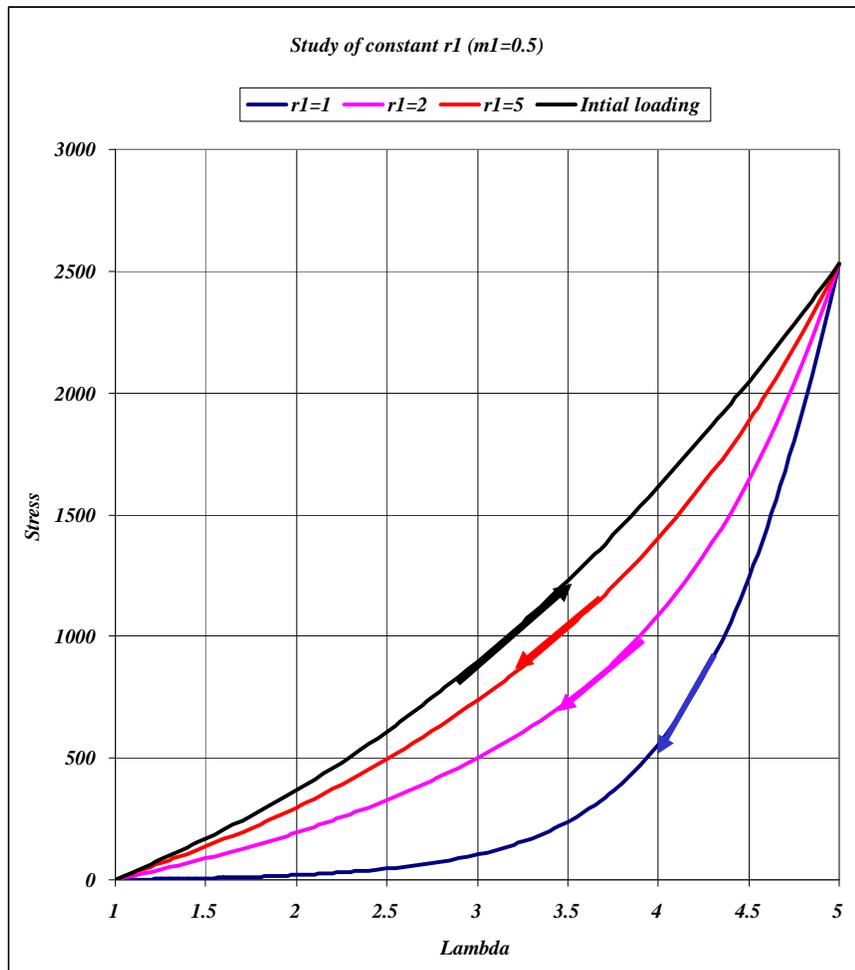


Figure 2. The effect of  $r_1$  on unloading for uniaxial extension

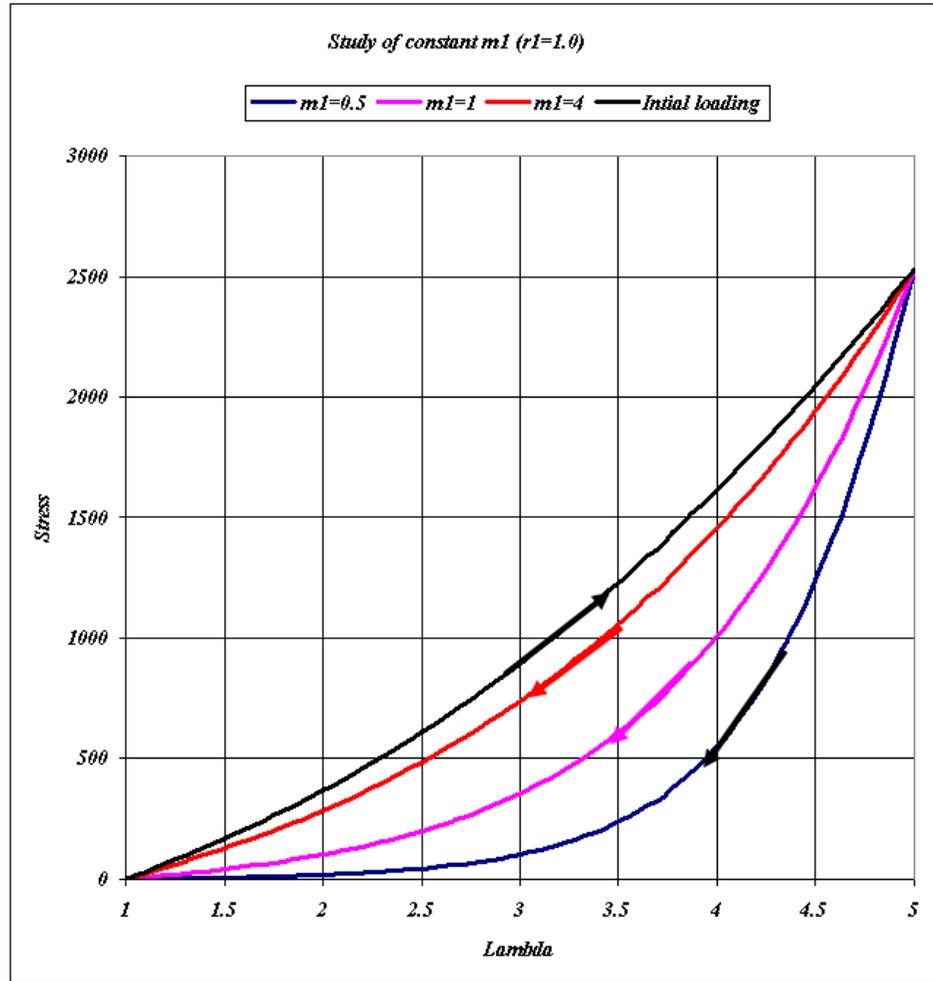


Figure 3. The effect of  $m_1$  on unloading for uniaxial extension

The effects of  $r_2$  and  $m_2$  for reloading are the same as  $r_1$  and  $m_1$ . With these four material constants the Mullins effect for most rubber-like materials can be modeled.

### Determining the damage constants

The material constants can be obtained from test data and the least-square error minimization method. The test datum for  $i^{\text{th}}$  stretch ratio on unloading or reloading is  $f[\lambda_1(i)]$ . Hence the error between test data and calculated value is

$$t_1[\lambda_1(i)] - f[\lambda_1(i)] \quad (11)$$

For  $m$  data points, the sum of the square of errors is

$$S = \sum_{i=1}^m \{t_1[\lambda_1(i)] - f[\lambda_1(i)]\}^2 \quad (12)$$

By minimizing the sum of the squares of errors  $S$ ,  $r_1$ ,  $m_1$ ,  $r_2$  and  $m_2$  are determined.

The initial loading curve, shown in Figure 4, is in red. The material constants are  $C_1=75$ ,  $\alpha = 0$ . In Part One of this paper the test data are generated numerically, and  $\pm 5\%$  experimental error was built into the numerical data. The raw data for unloading and reloading are shown in dots in Figure 4. The best fit that shows the Mullins effect and the raw data is shown in Figure 5. The determined Mullins damage constants are:  
 $r_1 = 0.8$ ,  $m_1 = 1.0$ ,  $r_2 = 0.5$  and  $m_2 = 5$ .

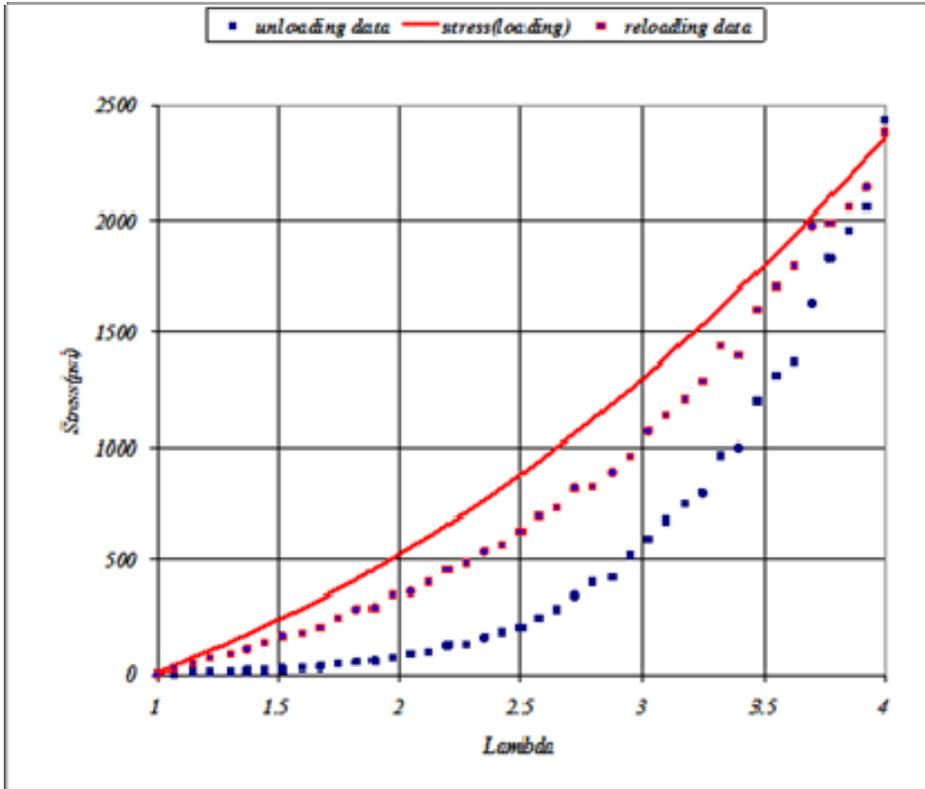


Figure 4. The stress-stretch curve for a neo-Hookean material and the numerical generated test data for unloading and reloading.

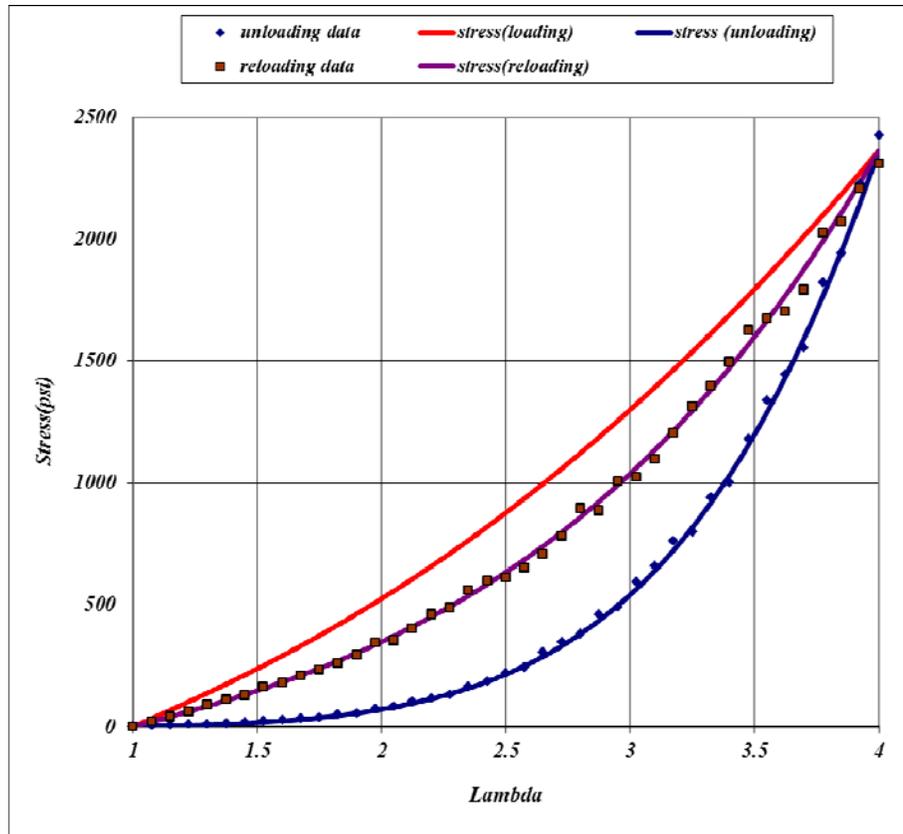


Figure 5. The best-fit unloading and reloading stress-stretch curves and the test data

### LS-DYNA implementation

The formulation presented in this paper applies to one-, two- and three-dimensional problems. For general three-dimensional problems the mathematical formulation has been implemented in LS-DYNA.

The result of a cube of  $0.5 \times 0.5 \times 0.5$  subjected to uniaxial extension is obtained from LS-DYNA. The displacement at one end and the stress in the cube are shown in Figure 6. The stress-displacement plot is shown in Figure 8. The same material constants, used in the analytical calculation, are used in the LS-DYNA calculation. The displacement ( $\Delta$ ) can be converted to the stretch ratio ( $\lambda$ ).  $\lambda = 1 + \Delta/L$ ; the undeformed length of the cube is  $L$ . The results shown in Figure 1 from the analytical calculations and Figure 7 from LS-DYNA are the same.

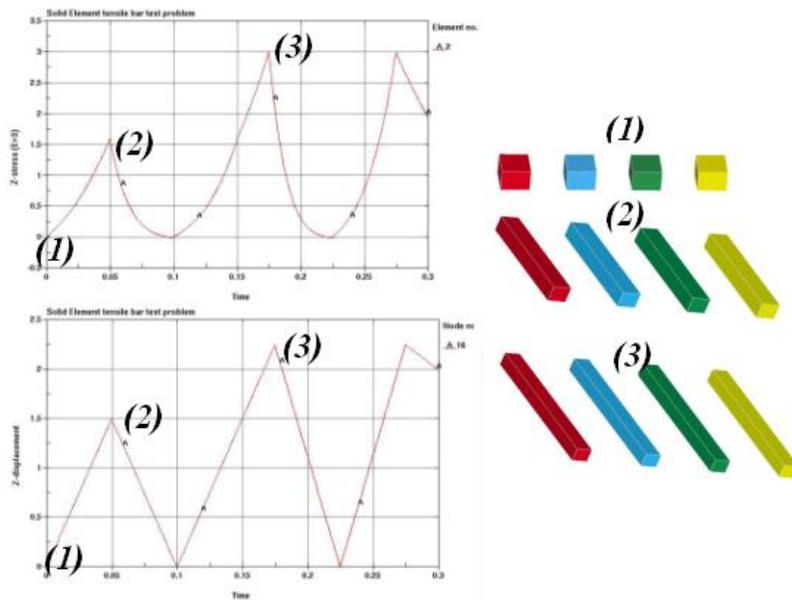


Figure 6. The results from LS-DYNA.

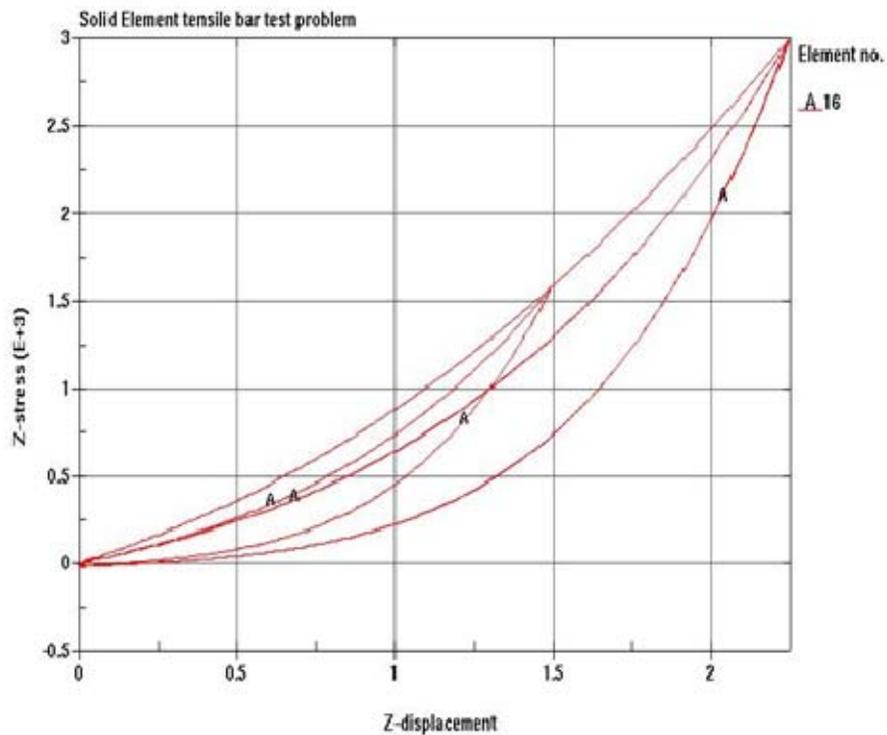


Figure 7. The stress-displacement plot from LS-DYNA.

### **Future work.**

In the second part of this paper, the Mullins effect on biaxial loading will be studied and a biaxial test will be performed. The material constants will be determined from the biaxial test data.

The formulations and applications can be extended to various rubbers with strain-energy density represented by various constitutive equations such as: neo-Hookean, Mooney, Ogden incompressible, and Ogden compressible materials. It can also be extended to viscoelastic materials for compressible and incompressible viscoelastic materials subjected to very large deformation.

We assumed that the Mullins damage function in this paper is represented by a hyperbolic tangent function; it can be changed to other functions if needed.

### **References**

1. W. W. Feng and J. O. Hallquist, "Numerical Modeling and Biaxial Tests for the Mullins Effect in Rubber, the 6<sup>th</sup> European LS-DYNA Conference, Gothenburg, Sweden, May 29-30, 2007.
2. W. W. Feng and J. O. Hallquist, "On Mooney-Rivlin Constants for Elastomers," the 12<sup>th</sup> International LS-DYNA Conference, Dearborn, Michigan, June 3-5, 2012.